# Autonomous Measurement-based Engine

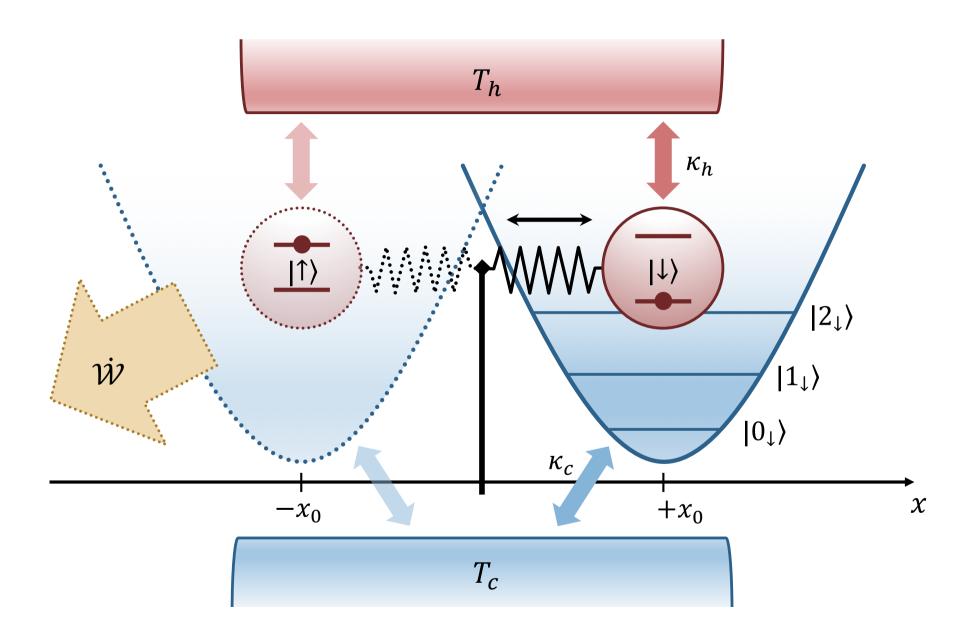
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#### Motivation

- Self-contained model
- Thermodynamic cost of measurement and erasure
- Origin of quantum heat



#### Model

Unitary component:

$$\hat{H} = \frac{\hbar\Omega}{2}\hat{\sigma}_z + \hbar\omega\left(\hat{a}^{\dagger}\hat{a} + \frac{1}{2}\right) + \hbar\omega x_0\hat{\sigma}_z\frac{\hat{a} + \hat{a}^{\dagger}}{\sqrt{2}}$$

- Measurement cost from work required to shift the pointer:
- Ideal pointer = small  $\omega$ , large  $x_0$

#### Model

#### • Unitary component:

$$\hat{H} = \frac{\hbar\Omega}{2}\hat{\sigma}_z + \hbar\omega\left(\hat{a}^{\dagger}\hat{a} + \frac{1}{2}\right) + \hbar\omega x_0\hat{\sigma}_z \frac{\hat{a} + \hat{a}^{\dagger}}{\sqrt{2}}$$

$$= \frac{\hbar\Omega}{2}\hat{\sigma}_z + \hbar\omega\hat{b}^{\dagger}\hat{b} + const$$

$$\hat{b} = \hat{a} + \frac{x_0}{\sqrt{2}}\hat{\sigma}_z$$

#### Model

Cold bath:

$$\mathcal{L}_c \rho = \kappa_c (\bar{n}_c + 1) \mathcal{D}[\hat{b}] \rho + \kappa_c \bar{n}_c \mathcal{D}[\hat{b}^{\dagger}] \rho$$

Hot bath:

$$\mathcal{L}_h \rho = \dots$$

 $\rightarrow$  For clearly separated states (large  $x_0$ )

$$\rho_{\infty} = (1 - p_{\infty})|g\rangle\langle g| \otimes \hat{D}\rho_{g}\hat{D}^{\dagger} + p_{\infty}|e\rangle\langle e| \otimes \hat{D}^{\dagger}\rho_{e}\hat{D}$$

- ullet Continuous incoherent process at a rate  $\gamma$ 
  - Simplest scenario: dichotomic projective measurement given by P (left) and 1-P (right)
  - Conditional spin-flip if we measure left

$$\mathcal{L}_m \rho = \gamma \mathcal{D}[\hat{\sigma}_x \hat{P}] \rho + \gamma \mathcal{D}[\mathbf{1} - \hat{P}] \rho$$

• Steady state deviates more from  $ho_{\infty}$  for high measurement rates

• Two energy terms due to measurement channel:

Work extracted + shift pointer

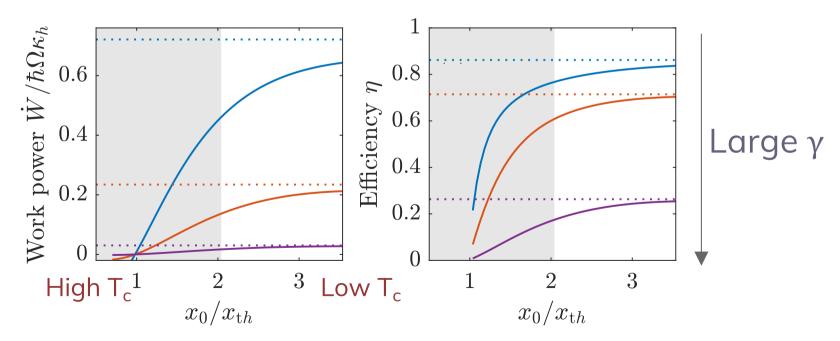
$$\operatorname{Tr}\left(\hat{H}\mathcal{L}_{m}\rho\right) = -\gamma\operatorname{Tr}\left(\hat{P}\rho_{\infty}\hat{P}\left[\hbar\Omega + 2\hbar\omega x_{0}\hat{x}\right]\hat{\sigma}_{z}\right)$$

$$+2\gamma\hbar\omega\operatorname{Tr}\left(\hat{b}^{\dagger}\hat{b}\mathcal{D}[\hat{P}]\rho_{\infty}\right)$$

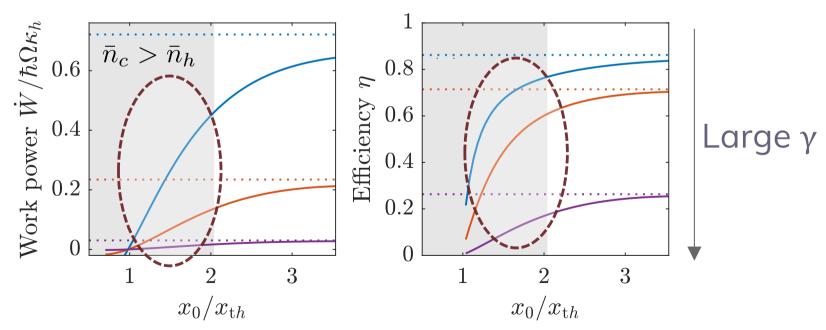
Backaction noise/dispersion

"Quantum heat"

• Ideal efficiency:  $\eta_{\rm max} pprox rac{1-2\omega x_0^2/\Omega}{1+2[1+(2ar{n}_h+2)\kappa_h/\gamma]\omega x_0^2/\Omega} < \eta_{\rm Otto}$ 



 But... larger operation window and no tradeoff between work and efficiency



## 2. Non-invasive probing

Apply position-dependent driving field:

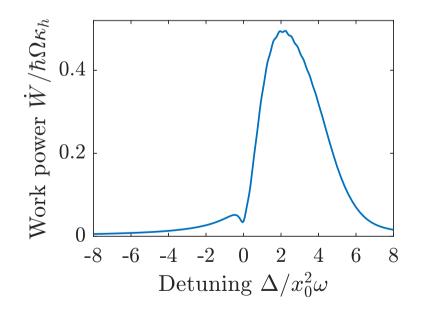
$$\hat{V}(t) = \hbar \zeta f(\hat{x}) e^{-i(\Omega - \Delta)t} |e\rangle \langle g| + h.c.$$

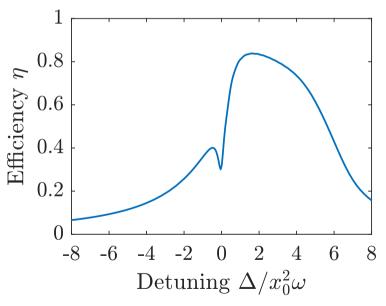
No ambiguity of work

$$\dot{W} = -\text{Tr}\left\{\rho_{\infty}(t)\partial_t \hat{V}(t)\right\}$$

## 2. Non-invasive probing

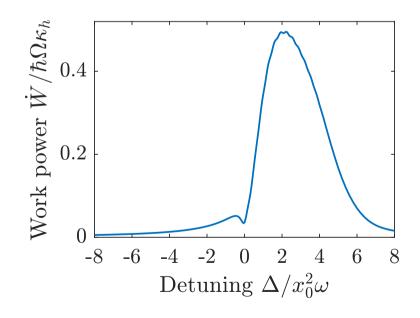
- Poorer than incoherent measurement
  - Finite feedback time and interaction strength

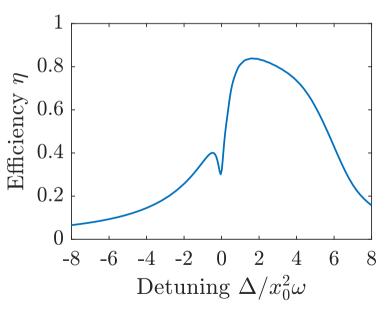




## 2. Non-invasive probing

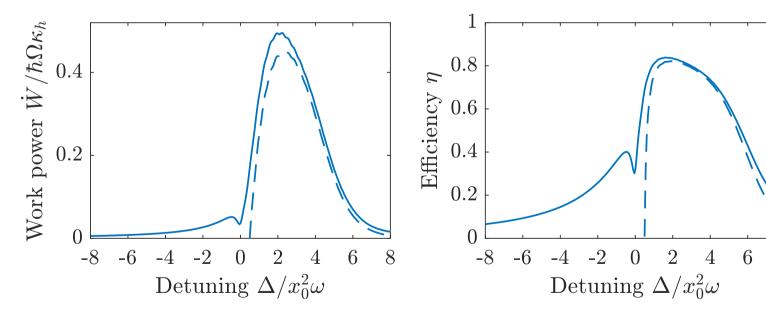
- High power & efficiencies for  $\Delta = 2\omega x_0^2$ 
  - = shift in qubit frequency when pointer is on left





#### **Alternative:**

 Probing through red-detuned fields that address qubit only when on left without position dependence



8

#### Conclusions

- (1) incoherent and (2) coherent
  - Also indirect measurement through frequency
- Simultaneous high work power and efficiencies
- Bigger operation window than Otto