



# Fast state preparation in trapped ion atomic clocks: **A fluctuation theorem perspective in the strong-coupling, non-Markovian regime**

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Braunschweig

Theory: PRB 94, 214308 (2016).  
Experiment: PRA 98, 023424 (2018).



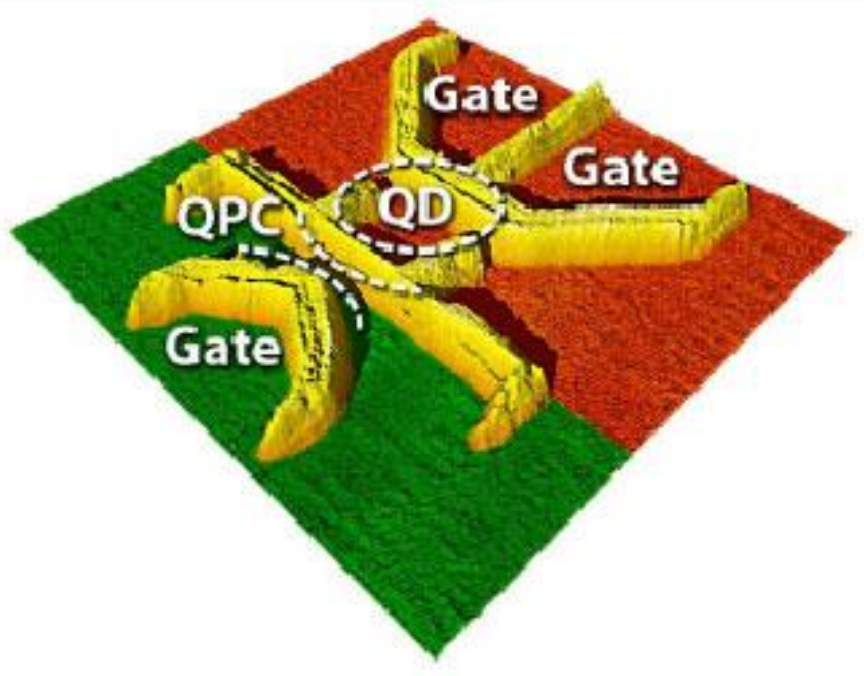
Physikalisch-Technische Bundesanstalt  
Braunschweig und Berlin

Quantum Thermodynamics Conference QTD2019  
Espoo, 26.06.2019



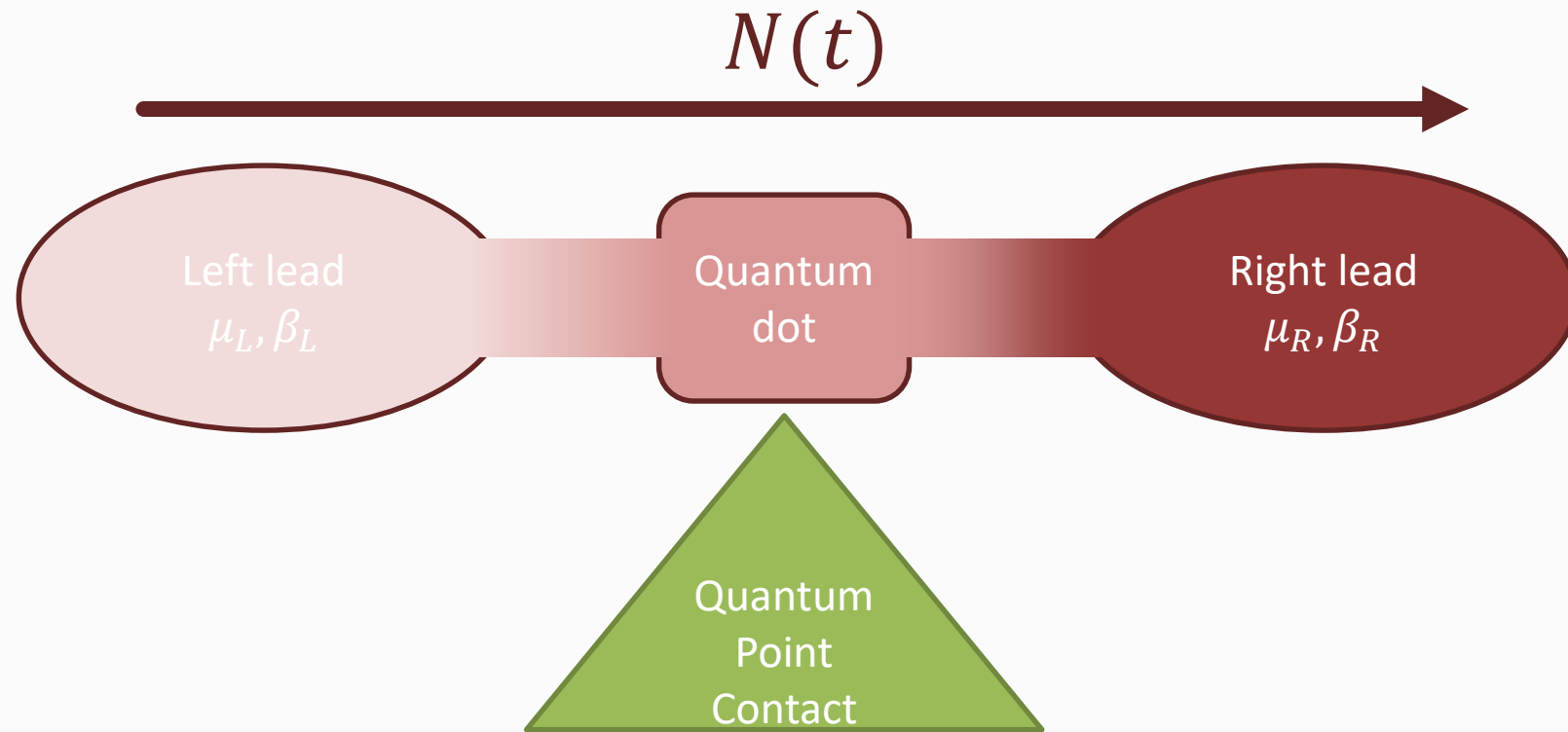
# Full Counting Statistics

Electronic transport through a quantum dot



R. Haug, Hannover Univ.  
J. Pekkola, Aalto Univ.

...



Cumulant generating function

$$C(\chi, t) = \ln \sum_n e^{i\chi N(t)} p[N(t)]$$

# Symmetries of Cumulant Generating Function



Shift of the counting field

Strong coupling to non-Markovian baths

Redefinition of system-environment boundaries

**Polaron transformation**

L. Nicolin and D. Segal, *PRB* 84, 161414(R) (2011).  
 L. Nicolin and D. Segal, *JCP* 135, 164106 (2011). ✓

**Consider environment part of the system**

M. Campisi, P. Talkner and P. Hänggi, *PRL* 102, 210401 (2009). ✓

**Chain Mappings / Reaction coordinate**

R. Rosenbach, J. Cerrillo, S. Huelga, J. Cao, M. Plenio, *NJP* 18, 23035 (2016).  
 S. Restrepo, J. Cerrillo, P. Strasberg, G. Schaller, *NJP* 20, 053063 (2017).  
 S. Restrepo, S. Böhling, J. Cerrillo, G. Schaller, accepted *PRB* (2019).

**System-environment correlations are essential for thermometry**  
 M. Buser, J. Cerrillo, G. Schaller, J. Cao, *PRA* 96, 062122 (2017).

$$\chi \rightarrow -\chi + i\beta$$

Forward /  
backward  
processes

- Fluctuation theorems

Free energy  
changes

- Jarzynski equality

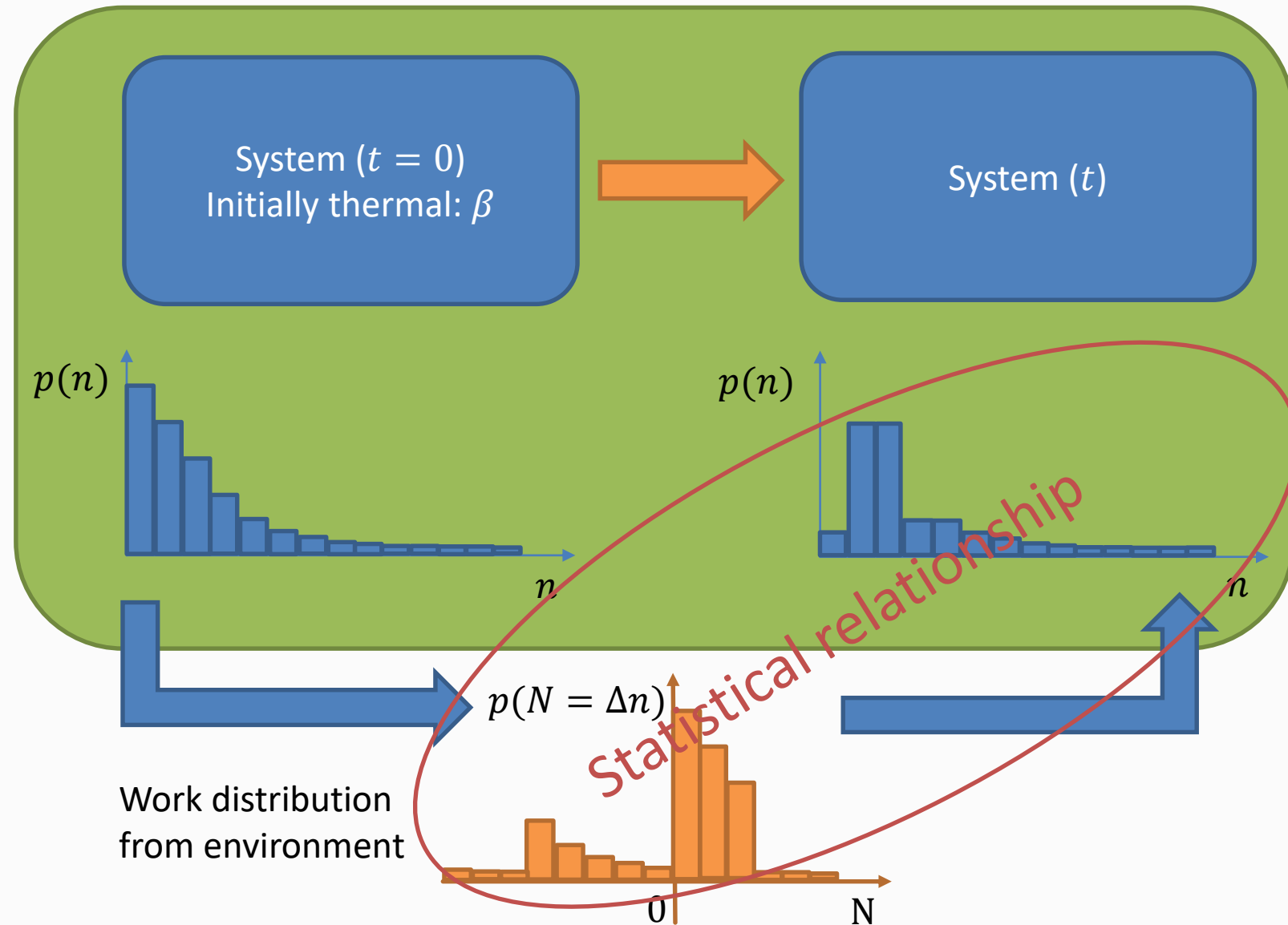
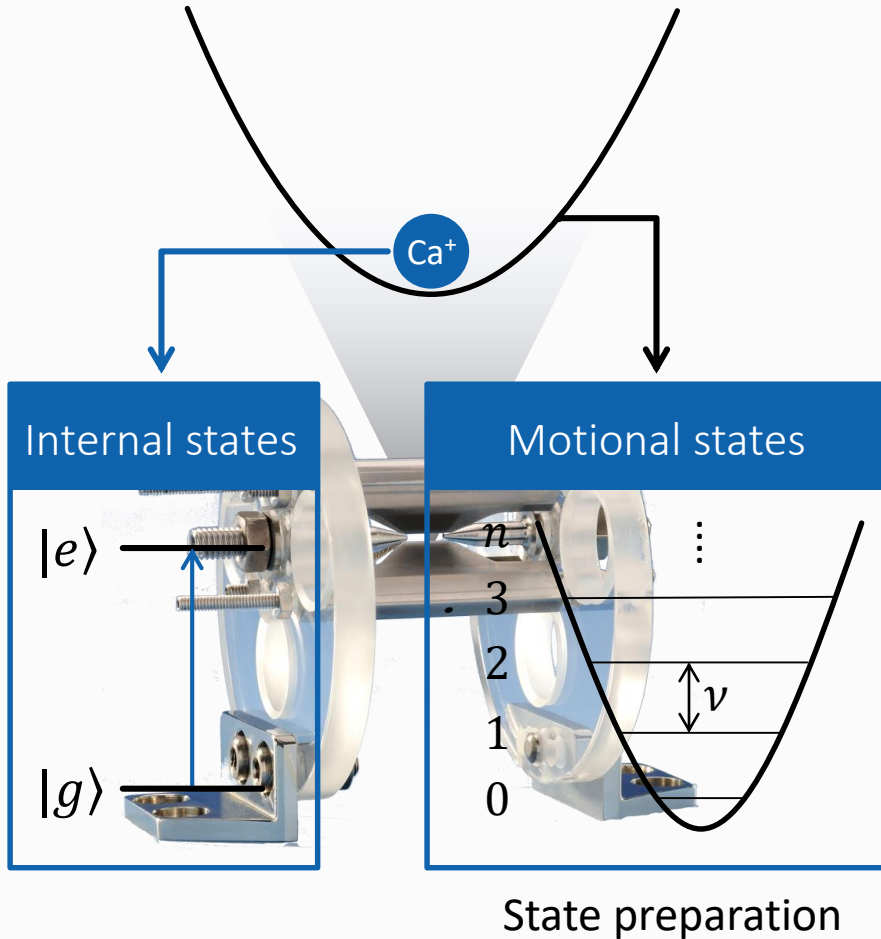
Detailed  
balance

Correlation Functions:  
Absorption and Emission  
 $A(t) \propto E(-t + i\beta)$

T H E R M O M E T R Y

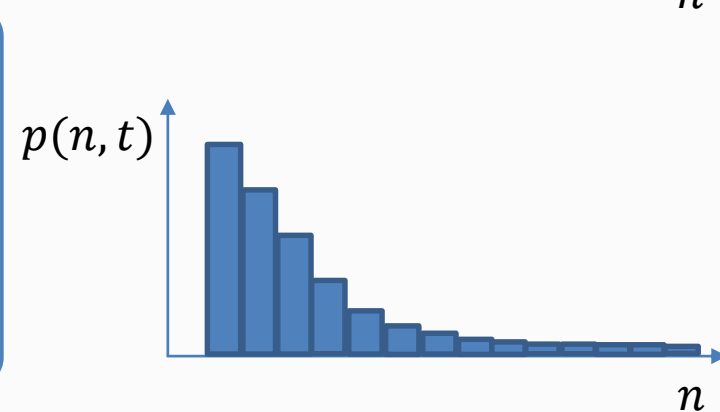
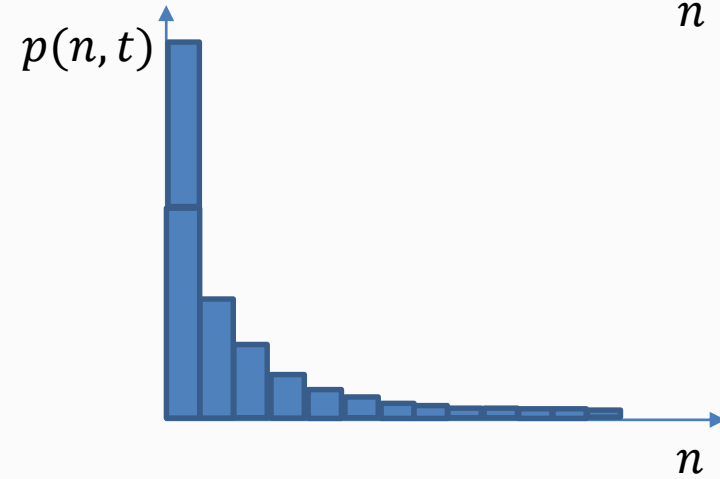
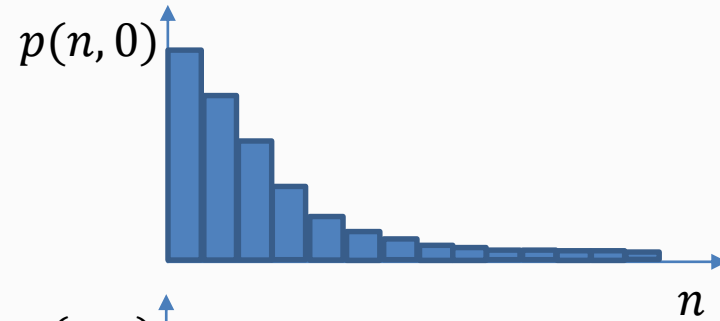
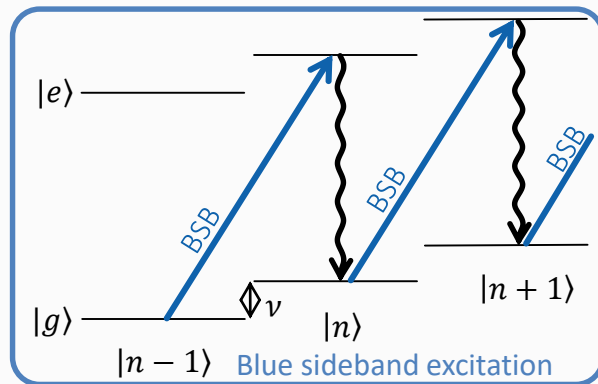
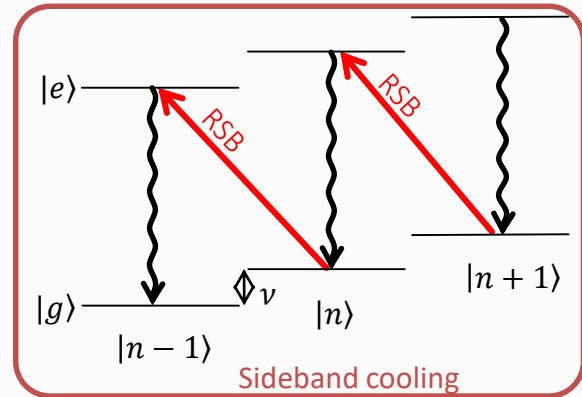
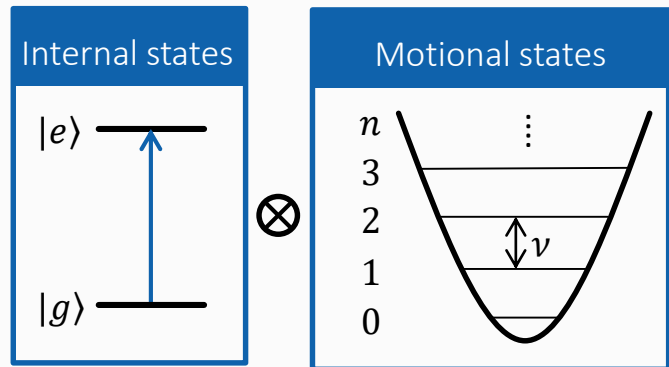


# Atomic clock preparation

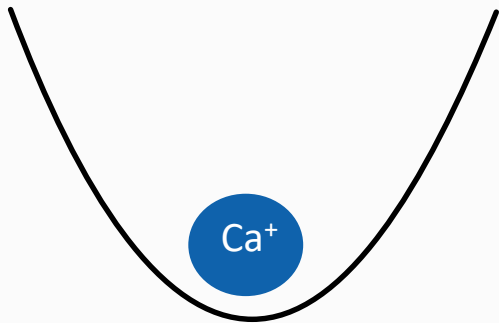
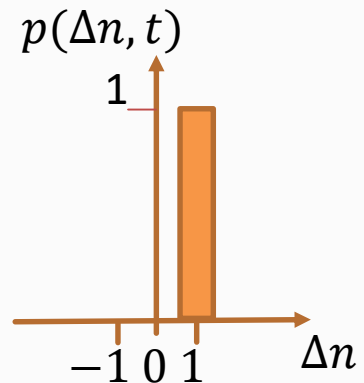
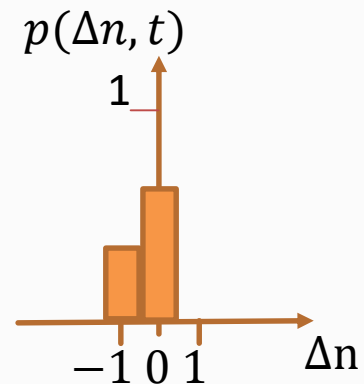




# Atomic clock preparation

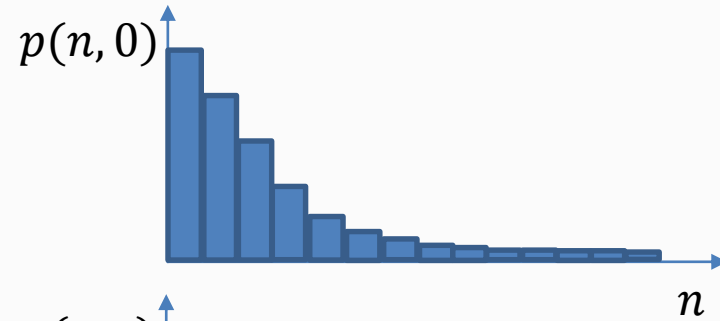
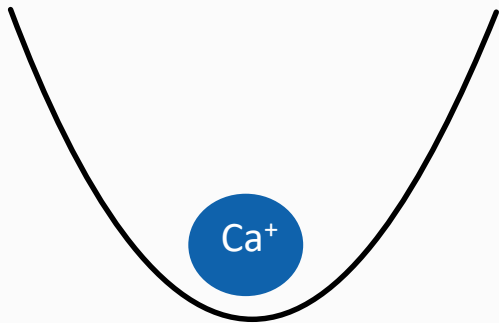
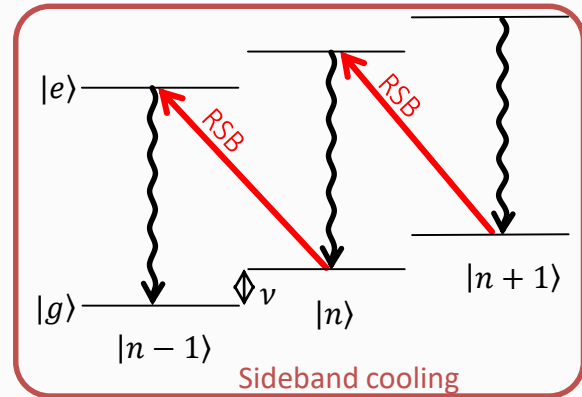
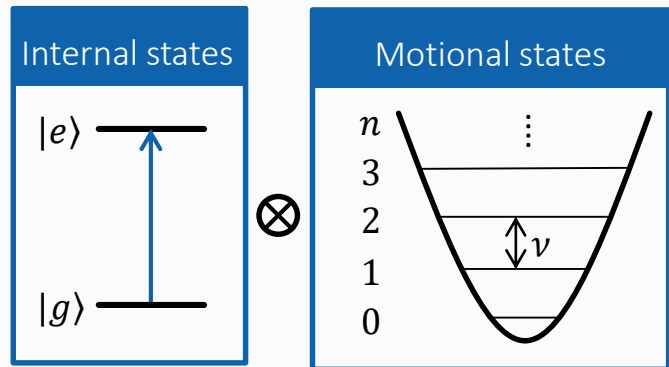


Transferred environmental energy (Work distribution)

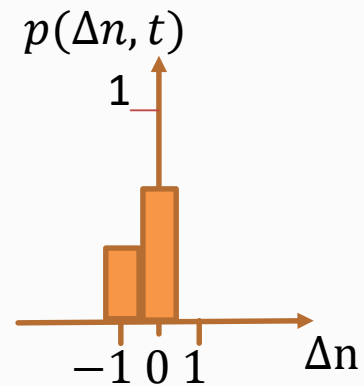
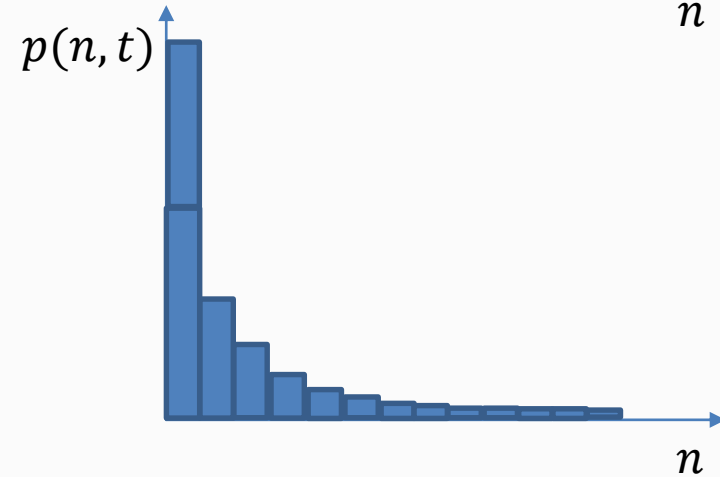




# Atomic clock preparation



Transferred environmental energy (Work distribution)



$$\frac{\int e^{i\chi n} p(n, t) dn}{\int e^{i\chi n} p(n, 0) dn} = \int e^{i\chi \Delta n} p(\Delta n, t) d\Delta n \Big|_{\beta=\beta-i\chi}$$

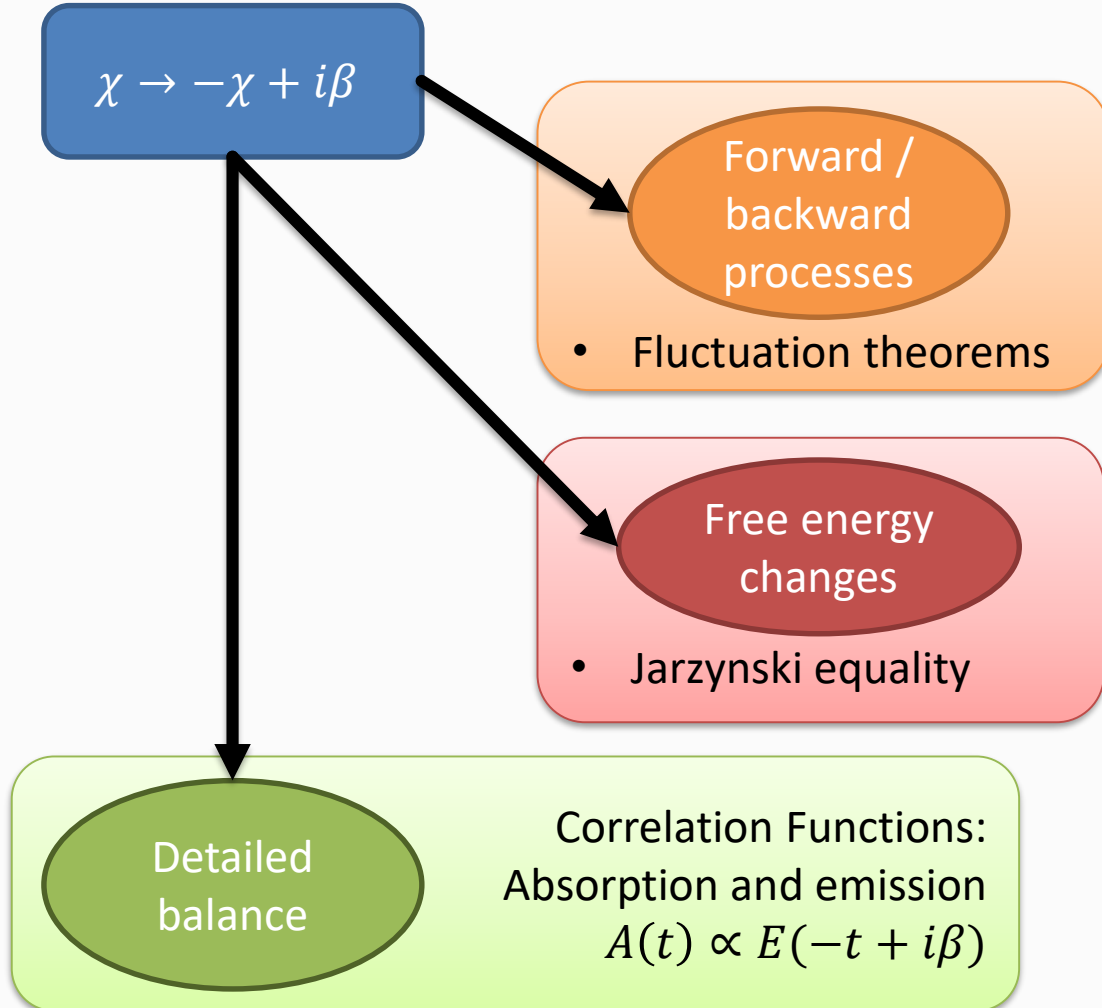
In terms of cumulant generating functions  $C(\chi, t, \beta)$ :

$$c(\chi, t, \beta) + F(\beta) - F(\beta - i\chi) = C(\chi, t, \beta - i\chi)$$



# Symmetries in Full Counting Statistics

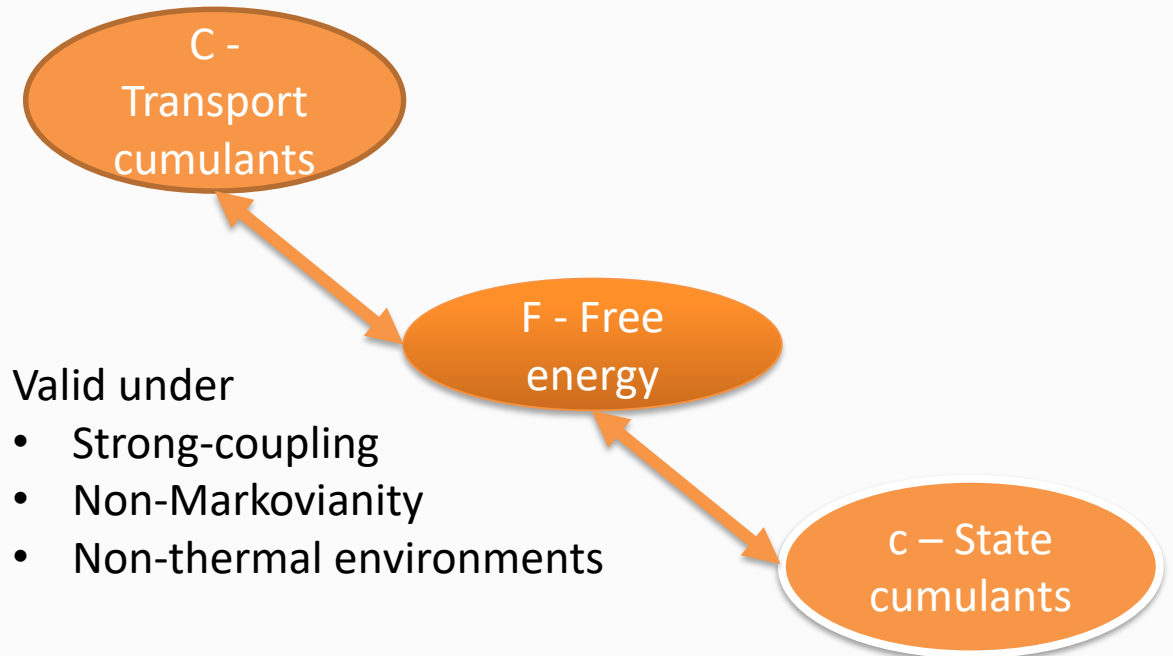
## Shift of the counting field



## Shift of the thermodynamic constraint

$$\beta \rightarrow \beta - i\chi$$

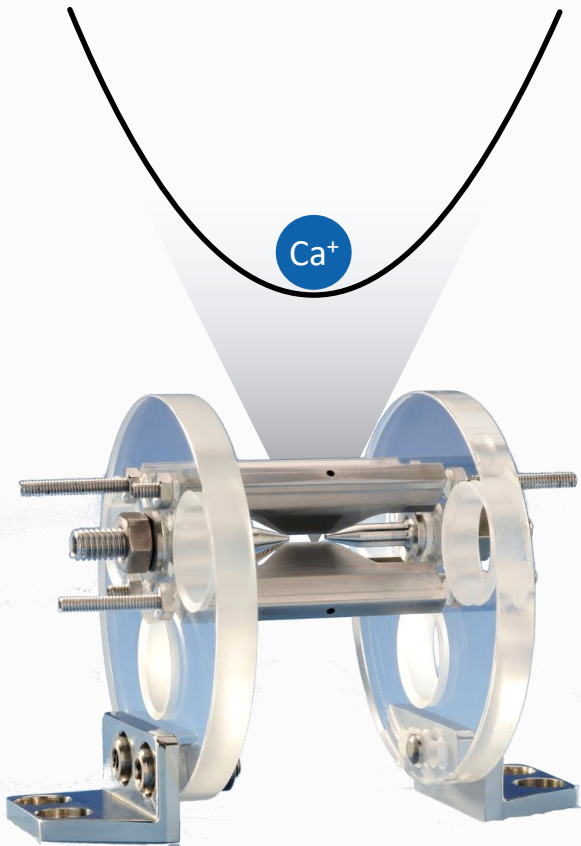
$$c(\chi, t, \beta) + F(\beta) - F(\beta - i\chi) = C(\chi, t, \beta - i\chi)$$





# Fast state preparation in trapped ion atomic clocks

- Attempt to go **beyond** the cooling speed limits of **sideband cooling**.
- Implementation of a double **electromagnetically-induced transparency** cooling scheme.
- The larger the **detuning** of the laser, the higher laser intensity can be used.



Rate as a figure of merit

$$R = \frac{\partial}{\partial \langle n \rangle} \langle \dot{n} \rangle$$

$$c(\chi, t, \beta) + F(\beta) - F(\beta - i\chi) = C(\chi, t, \beta - i\chi)$$

Fluctuation-dissipation like relation

$$R = \frac{-1}{2Cdt} [\Delta \langle \langle n^2(dt) \rangle \rangle - \langle \langle \Delta n^2(dt) \rangle \rangle]$$

Heat capacity of oscillator

Increase in state variance over dt

Variance of transferred phonons over dt

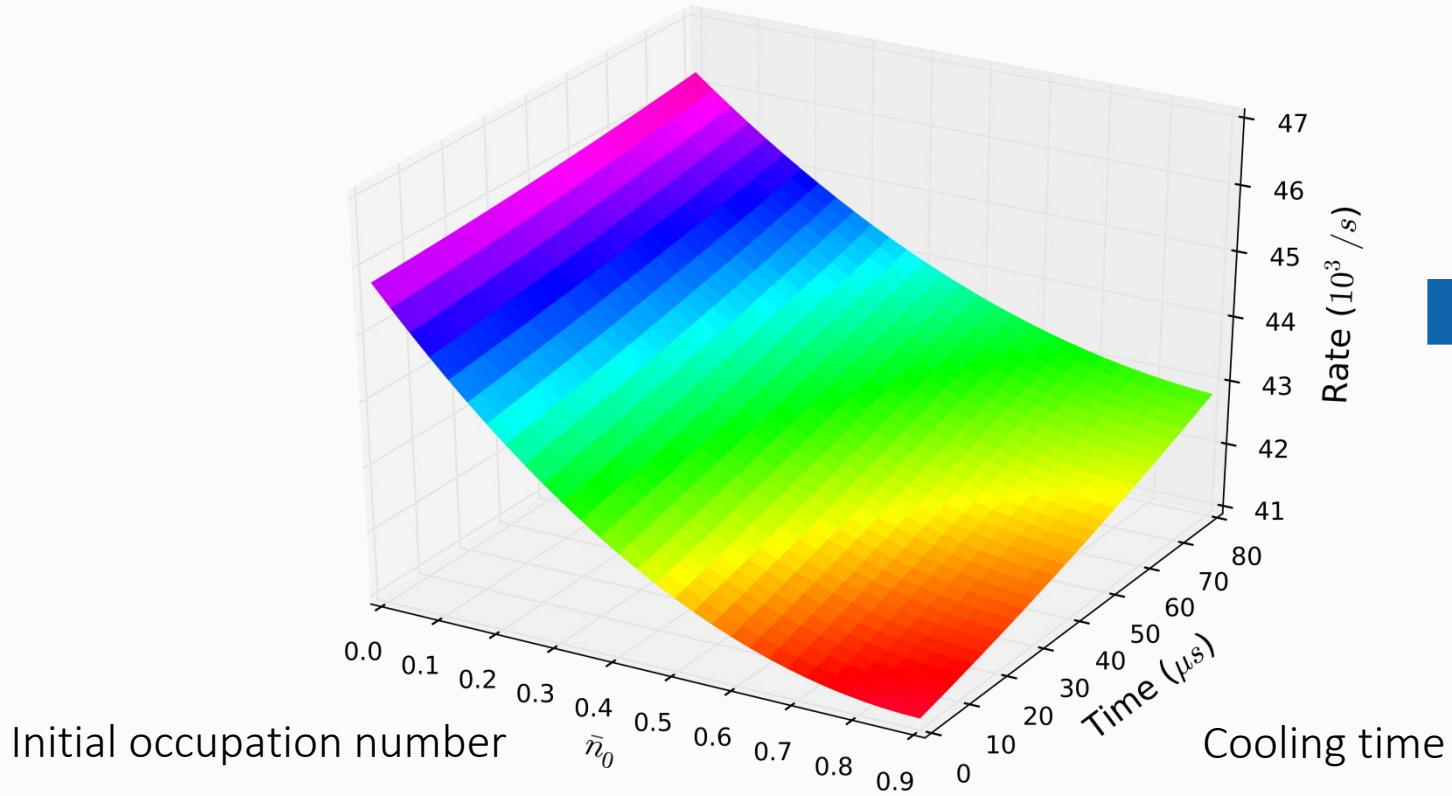
1. Weak coupling limit: Rate is independent of initial temperature
2. Increasing laser intensity: contribution to rate increase is detrimental and temperature dependent





# Full master prediction for cooling rate

Time dependent axial single EIT cooling rate  $R(t, n_0)$  at  $\Delta = 3\Gamma$

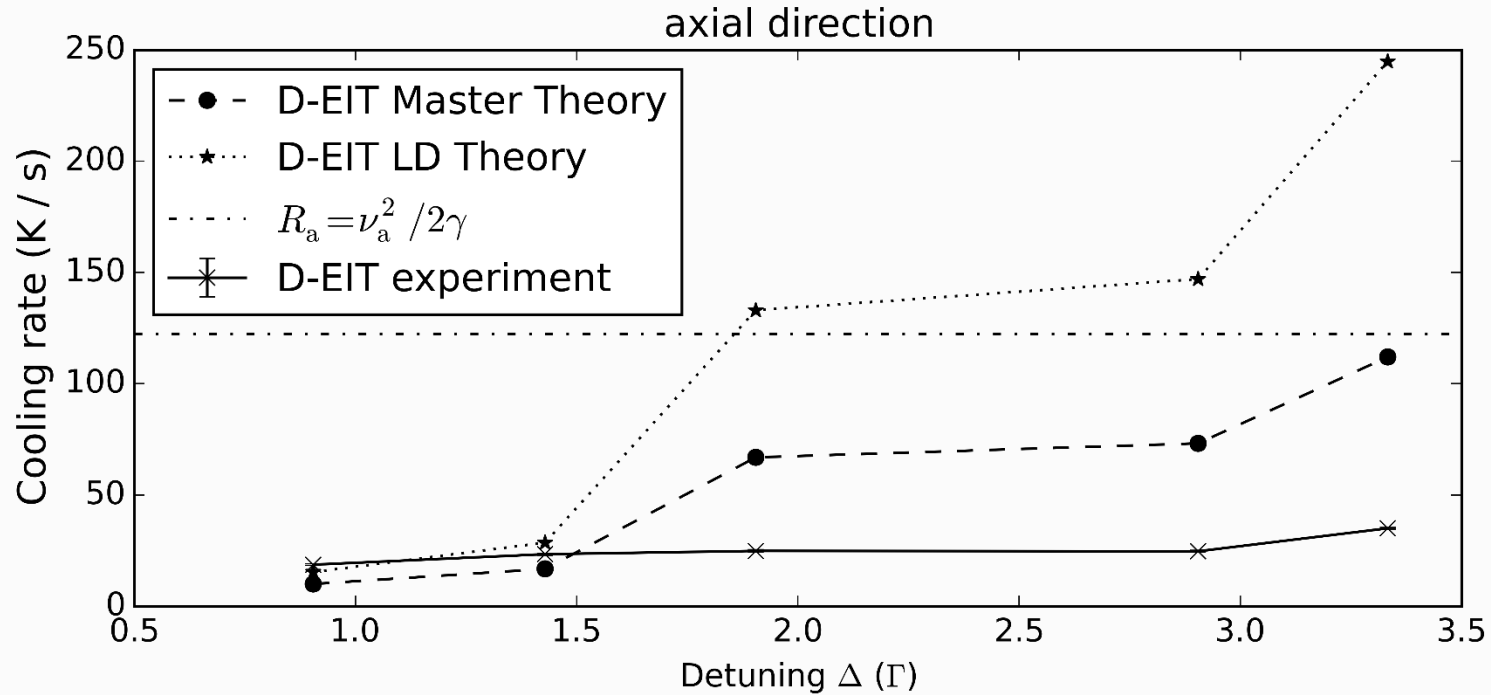


Lamb-Dicke theory:  $52.2 \times 10^3 s^{-1}$   
Measured:  $38.2 \times 10^3 s^{-1}$



# Full master prediction for cooling rate

## Lamb Dicke, Master equation and Experimental cooling rates



- Contrary to the prediction of the weak coupling theory, the cooling rate does not increase with detuning/laser intensity.
- There exists an upper bound to the cooling rate **imposed by the initial temperature** of the ion.



# Summary

$$c(\chi, t, \beta) + F(\beta) - F(\beta - i\chi) = C(\chi, t, \beta - i\chi)$$

- The preparation of a specific state involves an **exchange of an energy** distribution **with the environment**.
- Both distributions are non-trivially **related** through the expression above.
- The initial temperature of the system **imposes constraints** on state preparation
- These have been confirmed in **experiment** on atomic clock preparation.

*J. Cerrillo, M. Buser and T. Brandes, PRB 94, 214308 (2016).*

*N. Scharnhorst, J. Cerrillo et al., PRA 98, 023424 (2018).*